

# Dubin 4.2.2 Temperature Oscillations

11-15-16

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**Initialization:** Be sure the files *NTGStylesheet2.nb* and *NTGUtilityFunctions.m* are in the same directory as that from which this notebook was loaded. Then execute the cell immediately below by mousing left on the cell bar to the right of that cell and then typing “shift” + “enter”. Respond “Yes” in response to the query to evaluate initialization cells.

```
In[1]:= SetDirectory[NotebookDirectory[]];
(* set directory where source files are located *)
SetOptions[EvaluationNotebook[], (* load the StyleSheet *)
  StyleDefinitions -> Get["NTGStylesheet2.nb"]];
Get["NTGUtilityFunctions.m"]; (* Load utilities package *)
```

## Purpose

I solve a heat equation example from Chapter 4.2.2 of *Numerical and Analytical Methods for Scientists and Engineers, Using Mathematica*, Daniel Dubin. This example treats a time-oscillatory boundary condition.

## Solution

I construct an Association that encapsulates information about this problem, and then apply the function *DSolveHeatEquation* that attempts to solve this problem directly using the Mathematica function *DSolve*.

```
In[6]:= A1 =
Module[{description, pde, bcL, bcR, ic, eqns, depVar,
assumptions, substitutions, simplifications, names, values},
description = "Dubin 4.2.2 Temperature oscillations\nHomogeneous heat
equation, inhomogeneous time-dependentDirichlet boundary condition";
pde = D[T[x, t], t] - \[Chi] D[T[x, t], {x, 2}] == 0;
bcL = T[0, t] == T0; (* inhomogeneous Dirichlet *)
bcR = T[0, t] == T0 + TA Sin[\[Omega] t]; (* insulated, homogeneous von Neumann *)
ic = T[x, 0] == T0;
eqns = {pde, bcL, bcR, ic};
depVar = T[x, t];
assumptions = {L > 0, \[Chi] > 0};
substitutions = {K[1] \[Rule] n};
simplifications = {n \[Element] Integers};
values = {description, pde, bcL, bcR, ic,
eqns, depVar, assumptions, substitutions, simplifications};
names = {"description", "pde", "bcL", "bcR", "ic", "eqns",
"depVar", "assumptions", "substitutions", "simplifications"};
AssociationThread[names, values]];

Module[{soln, G},
soln = DSolveHeatEquation[A1];
AppendTo[A1, "soln" \[Rule] soln];
Print@ShowPDESetup[A1];
A1["soln"]]
```

**Dubin 4.2.2 Temperature oscillations**  
**Homogeneous heat equation, inhomogeneous time-dependentDirichlet boundary condition**

$$\begin{array}{ccc} T(0, t) = T_0 & \frac{\partial T(x, t)}{\partial t} - \chi \frac{\partial^2 T(x, t)}{\partial x^2} = 0 & T(0, t) = TA \sin(\omega t) + T_0 \\ \hline & T(x, 0) = T_0 & \end{array}$$

```
Out[7]= T^(0,1) [x, t] == \[Chi] T^(2,0) [x, t]
```

This problem is not immediately handled by DSolve.

The solution strategy is to write  $T(x,t)$  as the sum of a term that explicitly satisfies the boundary conditions  $T_{bc}(x)$  and a term that has homogeneous boundary conditions  $T_h(x,t)$ .

```
In[8]:= w[1] = T[x, t] == Tbc[x, t] + Th[x, t]
Out[8]= T[x, t] == Tbc[x, t] + Th[x, t]
```

A suitable choice for  $T_{bc}$  is

```
In[9]:= w[2] = Tbc[x, t] == T0 (1 - x/L) + x/L (T0 + TA Sin[ωt]) // Expand
Out[9]= Tbc[x, t] == T0 + x Sin[t ω] TA
L
```

The transformed heat equation is

```
In[10]:= w[3] = A1["pde"] /. T → Function[{x, t}, Th[x, t] + T0 + x Sin[t ω] TA];
w[3] = MapEqn[(# - x ω Cos[t ω] TA)/L &, w[3]];
Out[11]= Th^(0,1)[x, t] - χ Th^(2,0)[x, t] == -x ω Cos[t ω] TA
L
```

The boundary conditions for  $T_h(x, t)$  are

```
In[12]:= w[4] = {Th[0, t] == 0, Th[L, t] == 0}
Out[12]= {Th[0, t] == 0, Th[L, t] == 0}
```

The new initial condition for  $T_h(x, t)$  is

```
In[13]:= w[4] =
Th[x, t] == T[x, t] - Tbc[x, t] /. t → 0 /. T[x, 0] → T0 /. (w[2] /. t → 0 // ER)
Out[13]= Th[x, 0] == 0
```

A solution technique for such problems is to represent the source term and initial condition in terms of expansions of the eigenfunctions of the homogeneous equation for  $T_h(x, t)$ . Separation of variables leads to

```
In[14]:= w[5] = (A1["pde"][[1]] /. T → Th) == 0 /. Th → Function[{x, t}, τ[t] ψ[x]];
w[5] = MapEqn[(#/τ[t] ψ[x]) &, w[5]] // Expand
Out[15]= τ'[t]/τ[t] - χ ψ''[x]/ψ[x] == 0
```

The separated equations are

```
In[16]:= w[6] = {w[5][[1, 1]] == -λ, w[5][[1, 2]] == λ}
Out[16]= {τ'[t]/τ[t] == -λ, -χ ψ''[x]/ψ[x] == λ}
```

The latter constitutes a Sturm-Liouville equation (see Appendix A)

In[17]:=  $w[7] = \left\{ \lambda_n \rightarrow \frac{n^2 \pi^2 \chi}{L^2}, \psi_n[x] \rightarrow \sin\left[\frac{x \sqrt{\lambda_n}}{\sqrt{\chi}}\right] \right\}$

Out[17]=  $\left\{ \lambda_n \rightarrow \frac{n^2 \pi^2 \chi}{L^2}, \psi_n[x] \rightarrow \sin\left[\frac{x \sqrt{\lambda_n}}{\sqrt{\chi}}\right] \right\}$

Then, the solution for the homogeneous equation for  $T_h(x,t)$  is

In[18]:=  $w[8] = T_h[x, t] == \sum_{n=1}^{\infty} \tau_n[t] \psi_n[x]$

Out[18]=  $T_h[x, t] == \sum_{n=1}^{\infty} \tau_n[t] \psi_n[x]$

Substitute this into the inhomogeneous equation

In[19]:=  $w[9] = w[3] /. T_h \rightarrow \text{Function}\{x, t\}, \sum_{n=1}^{\infty} \tau_n[t] \psi_n[x]$

Out[19]=  $\sum_{n=1}^{\infty} \psi_n[x] \tau_n'[t] - \chi \sum_{n=1}^{\infty} \tau_n[t] \psi_n''[x] == -\frac{x \omega \cos[t \omega] T_A}{L}$

In[20]:=  $w[10] = w[9] /. \psi_n''[x] \rightarrow -\frac{\lambda_n}{\chi} \psi_n[x]$

Out[20]=  $-\chi \sum_{n=1}^{\infty} -\frac{\lambda_n \tau_n[t] \psi_n[x]}{\chi} + \sum_{n=1}^{\infty} \psi_n[x] \tau_n'[t] == -\frac{x \omega \cos[t \omega] T_A}{L}$

Write the source term in terms of an eigenfunction expansion

In[21]:=  $w[11] = w[10] /. -\frac{x \omega \cos[t \omega] T_A}{L} \rightarrow \sum_{n=1}^{\infty} f_n[t] \psi_n[x]$

Out[21]=  $-\chi \sum_{n=1}^{\infty} -\frac{\lambda_n \tau_n[t] \psi_n[x]}{\chi} + \sum_{n=1}^{\infty} \psi_n[x] \tau_n'[t] == \sum_{n=1}^{\infty} f_n[t] \psi_n[x]$

For each n

In[22]:=  $w[12] = w[11] /. \text{Sum}[a_, b_] \rightarrow a;$   
 $w[12] = \text{MapEqn}[(\# / \psi_n[x]) \&, w[12]] // \text{Expand}$

Out[23]=  $\lambda_n \tau_n[t] + \tau_n'[t] == f_n[t]$

The explicit  $f_n[t]$  are given by

```
In[24]:= w[13] = f_n[t] == Integrate[(x ω Cos[t ω] T_A Sin[n π x]/L), {x, 0, L}]/Integrate[Sin[n π x]^2, {x, 0, L}] // Refine[#, n ∈ Integers] &
Out[24]= f_n[t] == 2 (-1)^n ω Cos[t ω] T_A/(n π)
```

Thus

```
In[25]:= w[14] = w[12] /. (w[13] // ER)
Out[25]= λ_n τ_n[t] + τ_n'[t] == 2 (-1)^n ω Cos[t ω] T_A/(n π)
```

```
In[26]:= w[15] = DSolve[w[14], τ_n[t], t][[1, 1]] /. C[1] → A_n
Out[26]= τ_n[t] → e^{-t λ_n} A_n + (2 (-1)^n ω T_A (ω Sin[t ω] + Cos[t ω] λ_n)) / (n π (ω^2 + λ_n^2))
```

The value of the constant  $A_n$  are determined by the initial condition

```
In[27]:= w[16] = T_h[x, t] == Sum[τ_n[t] ψ_n[x], {n, 1, ∞}] /. w[15]
Out[27]= T_h[x, t] == Sum[e^{-t λ_n} A_n + (2 (-1)^n ω T_A (ω Sin[t ω] + Cos[t ω] λ_n)) / (n π (ω^2 + λ_n^2)) ψ_n[x], {n, 1, ∞}]
```

```
In[28]:= w[17] = w[16] /. t → 0 /. T_h[x, 0] → 0
Out[28]= 0 == Sum[A_n + (2 (-1)^n ω T_A λ_n) / (n π (ω^2 + λ_n^2)), {n, 1, ∞}] ψ_n[x]
```

Thus, each term of the summand must be zero, and the explicit  $A_n$  are

```
In[29]:= w[18] = Solve[A_n + (2 (-1)^n ω T_A λ_n) / (n π (ω^2 + λ_n^2)) == 0, A_n][[1, 1]]
Out[29]= A_n → -2 (-1)^n ω T_A λ_n / (n π (ω^2 + λ_n^2))
```

So, the complete solution for the temporal term is

```
In[30]:= w[19] = w[15] /. w[18] // Simplify
Out[30]= τ_n[t] → (2 (-1)^n e^{-t λ_n} ω T_A (e^{t λ_n} ω Sin[t ω] + (-1 + e^{t λ_n} Cos[t ω]) λ_n)) / (n π (ω^2 + λ_n^2))
```

and

In[31]:=

$$w[20] = w[8] /. \text{Sum} \rightarrow \text{Inactive}[\text{Sum}] /. w[19] // . w[7]$$

Out[31]=

$$T_h[x, t] = \sum_{n=1}^{\infty} \frac{2 (-1)^n e^{-\frac{n^2 \pi^2 t \chi}{L^2}} \omega \sin\left[\frac{\pi x \sqrt{\frac{n^2 \chi}{L^2}}}{\sqrt{\chi}}\right] \left( \frac{n^2 \pi^2 \chi \left(-1 + e^{\frac{n^2 \pi^2 t \chi}{L^2}} \cos[t \omega]\right)}{L^2} + e^{\frac{n^2 \pi^2 t \chi}{L^2}} \omega \sin[t \omega] \right) T_A}{n \pi \left(\frac{n^4 \pi^4 \chi^2}{L^4} + \omega^2\right)}$$

Finally, the solution of the original equation is

In[32]:=

$$w[21] = w[1] /. (w[2] // \text{ER}) /. (w[20] // \text{ER})$$

Out[32]=

$$T[x, t] = T_0 + \frac{x \sin[t \omega] T_A}{L} + \sum_{n=1}^{\infty} \frac{2 (-1)^n e^{-\frac{n^2 \pi^2 t \chi}{L^2}} \omega \sin\left[\frac{\pi x \sqrt{\frac{n^2 \chi}{L^2}}}{\sqrt{\chi}}\right] \left( \frac{n^2 \pi^2 \chi \left(-1 + e^{\frac{n^2 \pi^2 t \chi}{L^2}} \cos[t \omega]\right)}{L^2} + e^{\frac{n^2 \pi^2 t \chi}{L^2}} \omega \sin[t \omega] \right) T_A}{n \pi \left(\frac{n^4 \pi^4 \chi^2}{L^4} + \omega^2\right)}$$

I eliminate subscripts in preparation for a numerical treatment

In[33]:=

$$w[21] /. T_0 \rightarrow T0 /. T_A \rightarrow TA /. \infty \rightarrow nMax$$

Out[33]=

$$T[x, t] = T0 + \frac{TA x \sin[t \omega]}{L} + \sum_{n=1}^{nMax} \frac{2 (-1)^n e^{-\frac{n^2 \pi^2 t \chi}{L^2}} TA \omega \sin\left[\frac{\pi x \sqrt{\frac{n^2 \chi}{L^2}}}{\sqrt{\chi}}\right] \left( \frac{n^2 \pi^2 \chi \left(-1 + e^{\frac{n^2 \pi^2 t \chi}{L^2}} \cos[t \omega]\right)}{L^2} + e^{\frac{n^2 \pi^2 t \chi}{L^2}} \omega \sin[t \omega] \right)}{n \pi \left(\frac{n^4 \pi^4 \chi^2}{L^4} + \omega^2\right)}$$

In[34]:=

$$\text{Clear}[TSoln];$$

$$TSoln[x_, t_, L_, \chi_, T0_, TA_, \omega_, nMax_] :=$$

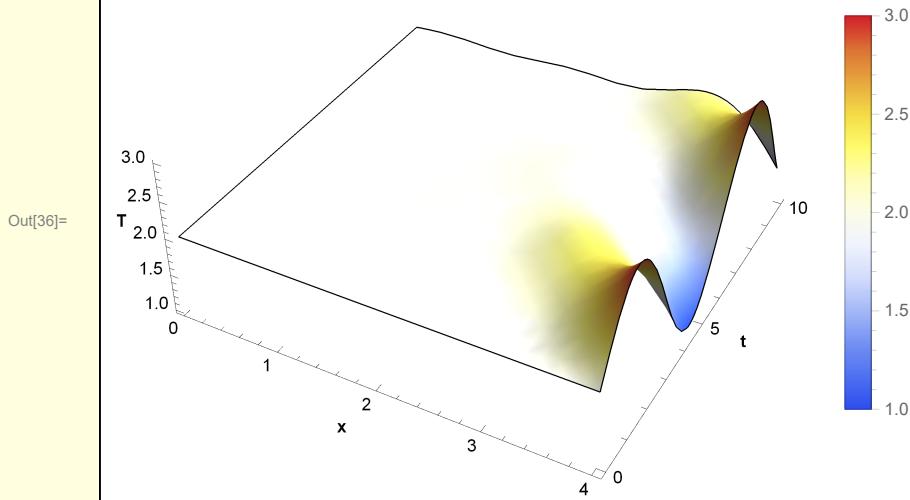
$$T0 + \frac{TA x \sin[t \omega]}{L} +$$

$$\sum_{n=1}^{nMax} \frac{2 (-1)^n e^{-\frac{n^2 \pi^2 t \chi}{L^2}} TA \omega \sin\left[\frac{\pi x \sqrt{\frac{n^2 \chi}{L^2}}}{\sqrt{\chi}}\right] \left( \frac{n^2 \pi^2 \chi \left(-1 + e^{\frac{n^2 \pi^2 t \chi}{L^2}} \cos[t \omega]\right)}{L^2} + e^{\frac{n^2 \pi^2 t \chi}{L^2}} \omega \sin[t \omega] \right)}{n \pi \left(\frac{n^4 \pi^4 \chi^2}{L^4} + \omega^2\right)} //$$

**Activate**

For the parameters specifies in Dubin's example (p287)

```
In[36]:= Module[{L = 4, x = 1/8, T0 = 2, TA = 1, ω = 1, nMax = 5},
Plot3D[TSoln[x, t, L, x, T0, TA, ω, nMax], {x, 0, 4}, {t, 0, 10}, PlotRange → All,
ColorFunction → "TemperatureMap", AxesLabel → {Stl["x"], Stl["t"], Stl["T"]},
Mesh → False, Boxed → False, PlotLegends → Automatic]]
```



I check this result by solving the same problem numerically

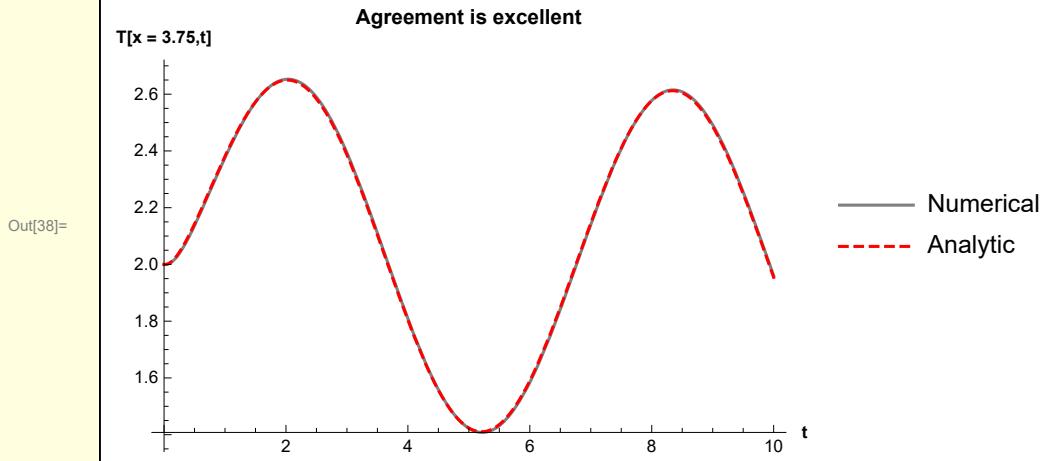
```
In[37]:= SOLNNUMERICAL = Module[{L = 4, x = 1/8, T0 = 2, TA = 1, ω = 1},
NDSolve[{D[T[x, t], t] == x D[T[x, t], {x, 2}], T[0, t] == T0,
T[L, t] == T0 + TA Sin[ω t], T[x, 0] == T0}, T, {x, 0, L}, {t, 0, 10}]]
```

```
Out[37]= T → InterpolatingFunction[ Domain: {{0., 4.}, {0., 10.}}]
```

Output: scalar

I compare TSoln against the numerical solutions as a function of time at the point x = 3.75.

```
In[38]:= Module[{L = 4, χ = 1/8, T0 = 2, TA = 1, ω = 1, nMax = 10, xTest = 3.75},
Plot[
{T[xTest, t] /. SOLNUMERICAL, TSoln[xTest, t, L, χ, T0, TA, ω, nMax]}, {t, 0, 10},
PlotStyle -> {Gray, Directive[Red, Dashed]}, PlotLegends -> {"Numerical", "Analytic"}, AxesLabel ->
{Stl["t"], Stl["T[x = 3.75,t]"]}, PlotLabel -> Stl["Agreement is excellent"]]
```



## Appendix Spatial eigenvalue problem

```
In[39]:= wA[1] = DSolve[-(χ ψ''[x]) == λ, ψ[x], x][[1, 1]] // RE
```

```
Out[39]= ψ[x] == C[1] Cos[x Sqrt[λ]/Sqrt[χ]] + C[2] Sin[x Sqrt[λ]/Sqrt[χ]]
```

```
In[40]:= wA[2] = wA[1] /. x → 0 /. ψ[0] → 0
```

```
Out[40]= 0 == C[1]
```

```
In[41]:= wA[3] = wA[1] /. Solve[wA[2], C[1]][[1]]
```

```
Out[41]= ψ[x] == C[2] Sin[x Sqrt[λ]/Sqrt[χ]]
```

The eigenvalues are given by

```
In[42]:= wA[4] = Solve[x Sqrt[λ] == n π, λ][[1, 1]] /. λ → λn
Out[42]= λn →  $\frac{n^2 \pi^2 \chi}{x^2}$ 
```

## Functions

```
In[3]:= Clear>ShowPDESetup];
ShowPDESetup[A_] := Module[{top = 1.0, right = 1.0,
  boundaries, labels, textInterior, textIC, textBCL, textBCR},
  boundaries = Line /@ {{{0, 0}, {right, 0}},
    {{0, 0}, {0, top}}, {{right, 0}, {right, top}}};
  labels = Text[PhysicsForm[A[[#][1]], #][2]] & /@
    {{{"pde", {right/2, top/2}}, {"ic", {right/2, 0.0}}},
     {"bcl", {0.0, top/2}}, {"bcR", {right, top/2}}};
  Graphics[{Directive[Black, Thick], boundaries, labels}, Axes → False,
  AspectRatio → 0.25, ImageSize → 500, PlotLabel → Stl[A["description"]]]]
```

```
In[4]:= Clear>DSolveHeatEquation];
DSolveHeatEquation[A_] :=
Module[{soln},
  soln =
  DSolve[A["eqns"], A["depVar"], {x, t}, Assumptions → A["assumptions"]][[1, 1]];
  soln = soln //. A["substitutions"];
  soln = Simplify[soln, A["simplifications"]];
  soln]
```